



# Fermionic Paired Superfluids in Rotating Trap

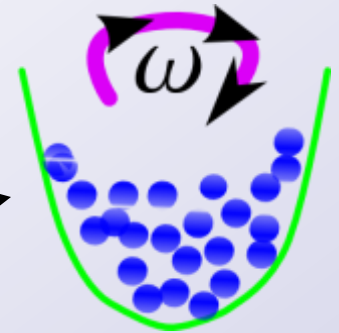
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# Outline

- Motivation:
  - Effect of rotation on degenerate Fermi gas across Feshbach resonance (Experiments with  ${}^6\text{Li}$  [Zweirlen et al. Nature, 2005] )
  - Fermions vs. Bosons:  
Additional mechanisms suppress  $T_c$  for fermions (Gorkov).
- Two-Channel Model In a Rotating Trap
  - Model
  - Perturbative Expansion
- Results
  - Non-Rotating Trap
  - Rotating Trap + Scaling
- Future Work

# Two Channel Model: Fermion + Boson in Rotating Trap



$$V(\mathbf{r}) = \frac{m}{2} (\Omega_{\perp}^2 (x^2 + y^2) + \Omega_z^2 z^2)$$

**In Rotating Frame:**  $\mathcal{H} = \mathcal{H}^{(F)} + \mathcal{H}^{(B)} + \mathcal{H}^{(I)}$

$$\mathcal{H}^{(F)} = \sum_{\sigma} \int d^3\mathbf{r} \psi_{\sigma}^{\dagger}(\mathbf{r}) \left[ -\frac{\hbar^2 \nabla_{\mathbf{r}}^2}{2m} + V(\mathbf{r}) - \mu - \omega L^z \right] \psi_{\sigma}(\mathbf{r})$$

$$\mathcal{H}^{(B)} = \int d^3\mathbf{r} \phi^{\dagger}(\mathbf{r}) \left[ -\frac{\hbar^2 \nabla_{\mathbf{r}}^2}{2(2m)} + 2V(\mathbf{r}) - 2\mu + \delta - \omega L^z \right] \phi(\mathbf{r})$$

$$\mathcal{H}^{(I)} = g \int d^3\mathbf{r} \left( \phi^{\dagger}(\mathbf{r}) \psi_{\uparrow}(\mathbf{r}) \psi_{\downarrow}(\mathbf{r}) + \psi_{\uparrow}^{\dagger}(\mathbf{r}) \psi_{\downarrow}^{\dagger}(\mathbf{r}) \phi(\mathbf{r}) \right)$$



# Procedure for calculation of $T_c$ (Perturbative expansion in $g^2$ )

- Imaginary time action
- Integrate out Fermions
- Expand out  $Z$  to second order in Bosonic fields

$$Z = \int \mathcal{D}\phi \mathcal{D}\psi e^{-S}$$

$$S = -T \text{Tr} \log \left[ \begin{pmatrix} -\mathcal{G}_0^F & g\phi \\ g\phi & -\mathcal{G}_0^B \end{pmatrix} \right] + \int d\tau d^3r \bar{\phi} \left( -\mathcal{G}_0^B \right)^{-1} \phi$$

■ given by Thouless criterion subject to constraint by particle number

$$S \approx -\log Z_0 + \int d\tau d^3r \bar{\phi}(\mathbf{r}) \left( -\mathcal{G}_0^B \right)^{-1} \phi(\mathbf{r}) - g^2 \int d\tau d^3r d^3r' \bar{\phi}(\mathbf{r}) \Pi(\mathbf{r}, \mathbf{r}') \phi(\mathbf{r}')$$

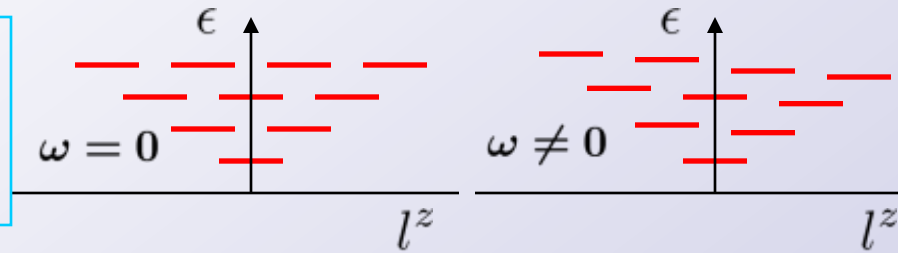
$$\Pi(\mathbf{r}, \mathbf{r}') = \mathcal{G}_0^F(\mathbf{r}, \mathbf{r}') \mathcal{G}_0^F(\mathbf{r}, \mathbf{r}')$$

$$\text{Min}_\lambda \left[ \left( \mathcal{G}_0^B \right)^{-1}(0) + g^2 \Pi(0) \right] = 0$$

$$N = N_0^F + 2N^B \text{ where } N^B = \frac{1}{\beta} \sum_n \frac{e^{i\Omega_n 0^+}}{\left( \mathcal{G}_0^B \right)^{-1}(i\Omega_n) + g^2 \Pi(i\Omega_n)}$$

## Using Harmonic Oscillator Basis

$$\begin{aligned}\epsilon_{l=\{L,l^z,l_0\}}^F &= \hbar\Omega_{\perp}(2L+1) + 2\hbar\omega l^z + \hbar\Omega_z l_0 - \mu \\ \epsilon_{l=\{L,l^z,l_0\}}^B &= \hbar\Omega_{\perp}(2L+1) + 2\hbar\omega l^z + \hbar\Omega_z l_0 - 2\mu + \delta\end{aligned}$$



## Thouless Criterion+ Particle Number

$$\textcircled{1} \quad \delta - 2\mu = \frac{g^2}{4\pi^{5/2} \bar{a}_{HO}^3} \sum_{l=\{L,l^z,l_0\}} \frac{e^{-\frac{2(l^z)^2}{L}}}{\sqrt{Ll^0}} \left[ \frac{\tanh(\beta\epsilon_l^F/2) + \tanh(\beta\epsilon_l^B/2)}{\epsilon_l^F + \epsilon_l^B} - \frac{2}{\epsilon_l^F + \epsilon_l^B + 2\mu} \right]$$

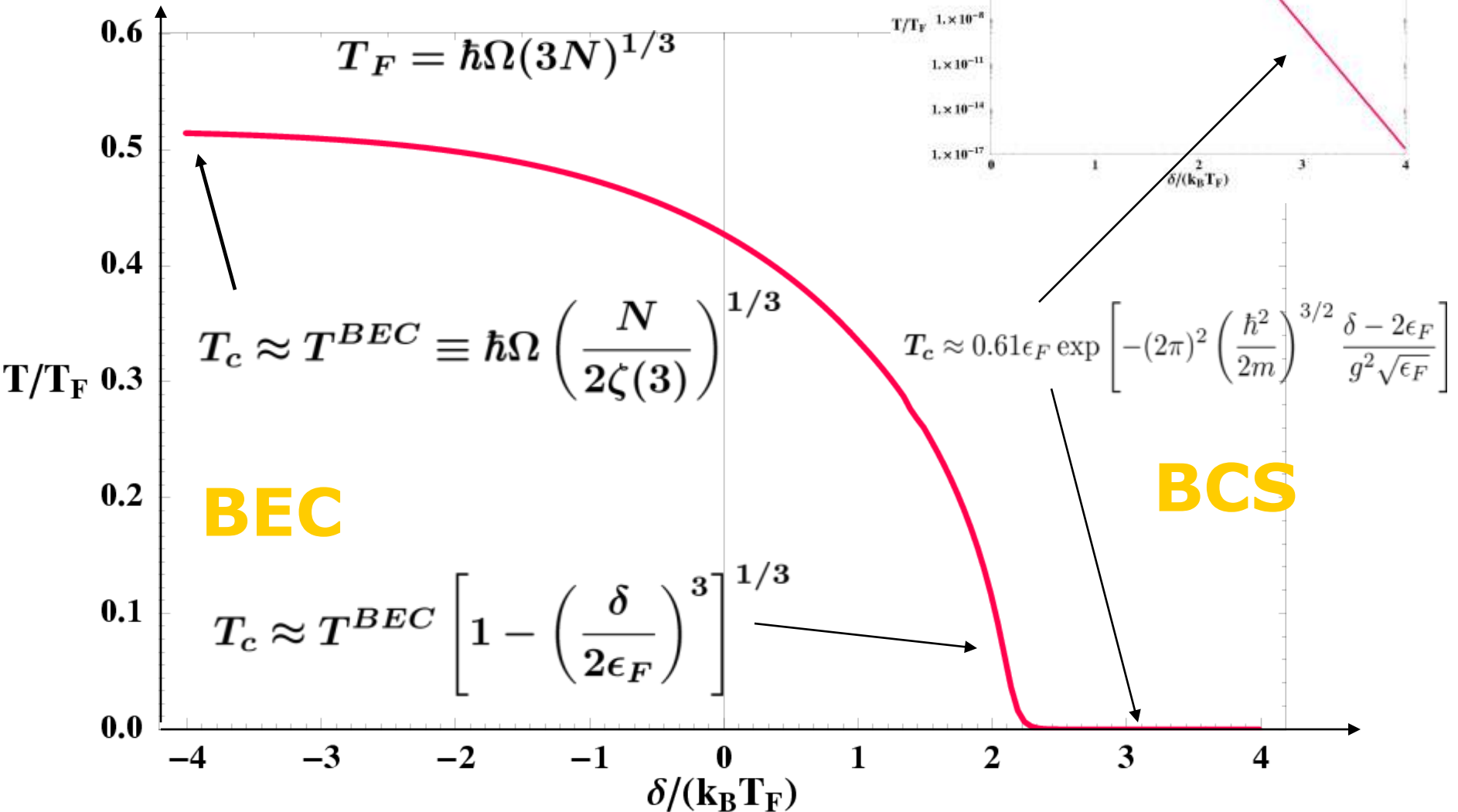
$$\textcircled{2} \quad N = \frac{2}{(\beta\hbar\tilde{\Omega})^3} \left( Li_3 \left[ e^{\beta(2\mu-\delta)} \right] - Li_3 \left[ -e^{\beta\mu} \right] \right) + \mathcal{O}(g^2) \text{ where } \tilde{\Omega} = \Omega \left[ 1 - \left( \frac{\omega}{\Omega_{\perp}} \right)^2 \right]^{1/3}$$

$\omega$  dependence

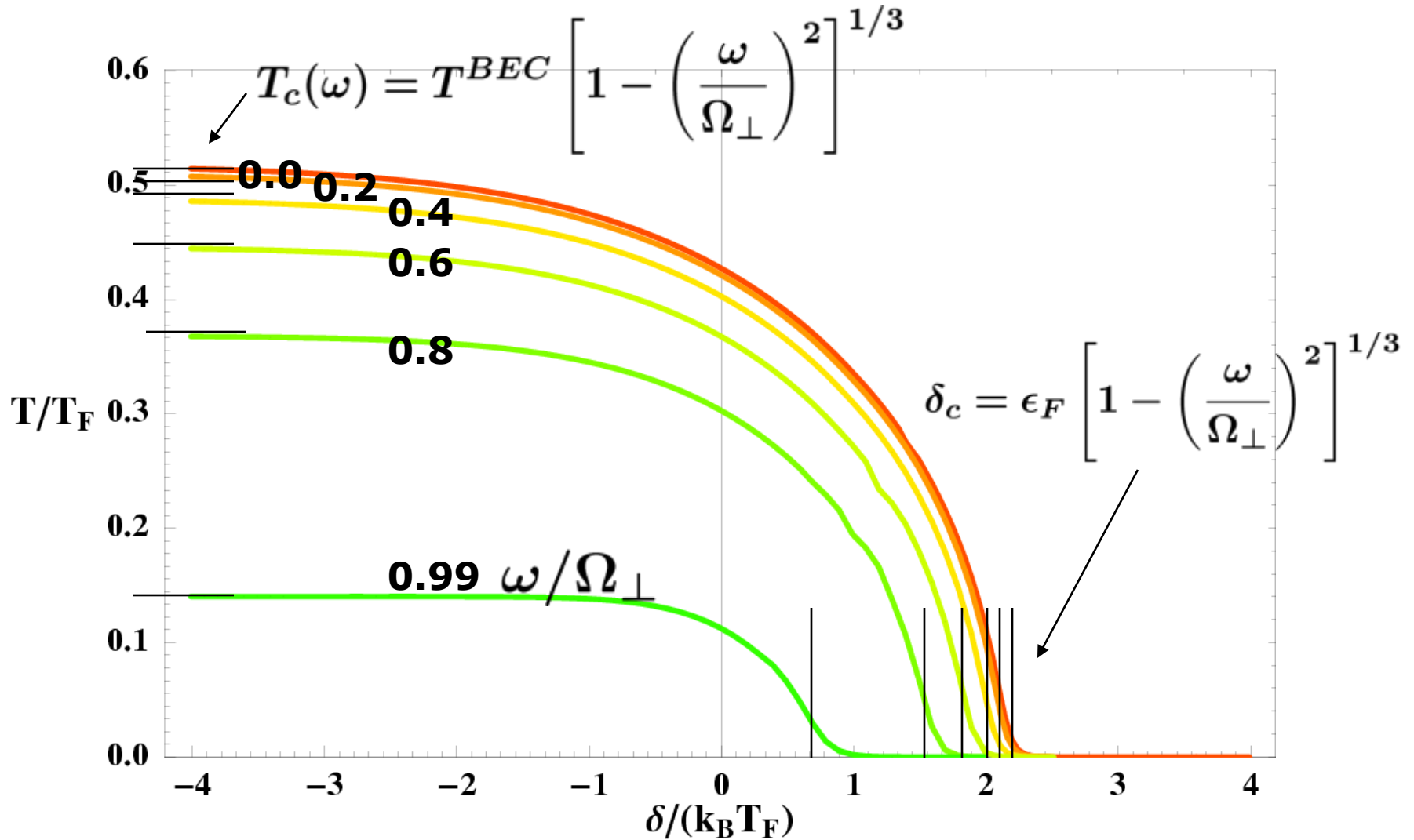
**Solve for  $(\mu, \beta)$**

- 1 Depairing effect: (Gorkov)**  
Rotation  $\iff$  Magnetic Fields in SC
- 2 Weakening of the trap:**  
lowering of the local density

# Result I: No Rotation ( $\omega=0$ )



# Result II: Rotating Trap





## Result III

- Scaling under rotation in Zeroth order in  $g$
- Further corrections cannot be obtained by rescaling interactions

Eg: BCS regime

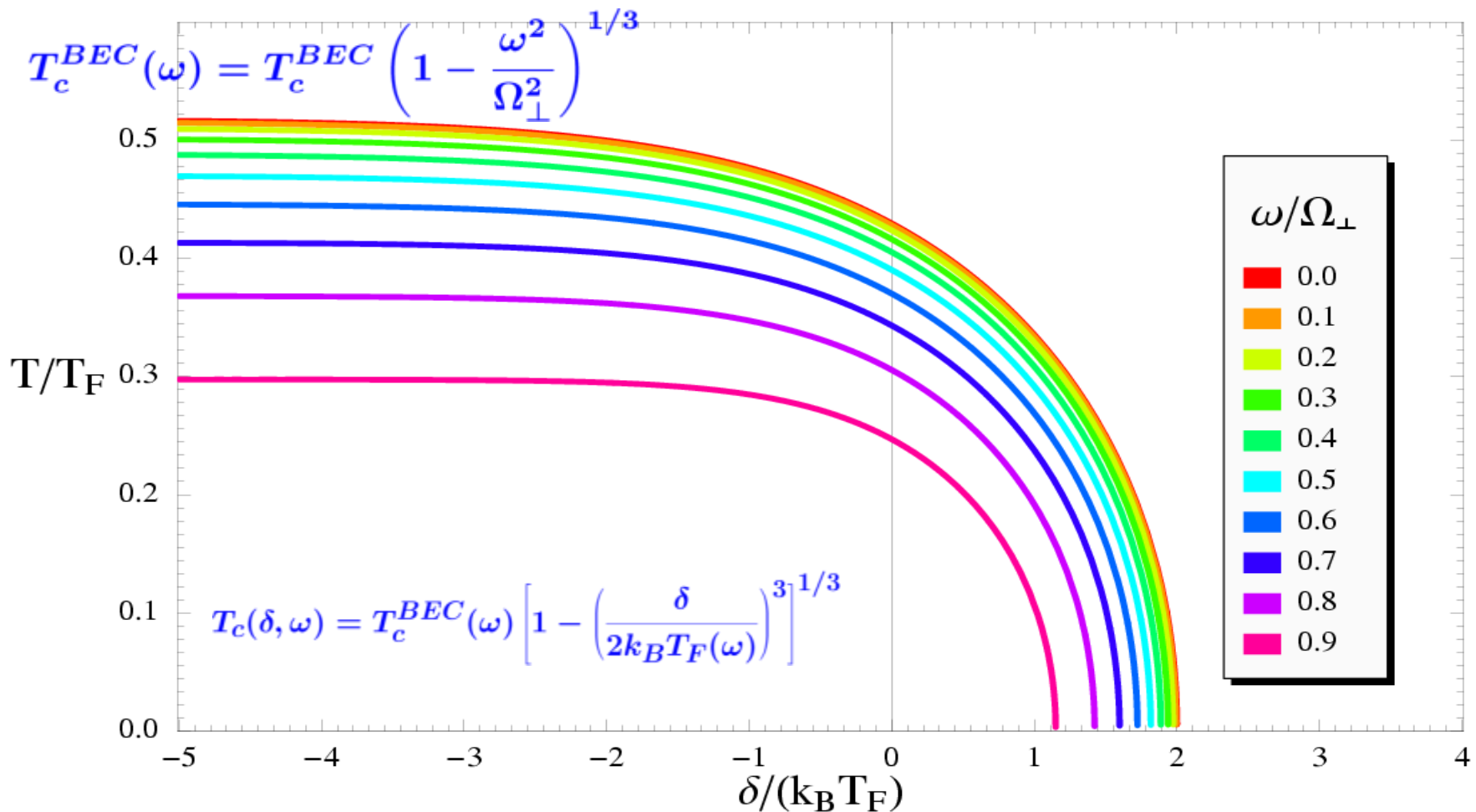
$$\delta \rightarrow \delta \left[ 1 - \left( \frac{\omega}{\Omega_{\perp}} \right)^2 \right]^{1/3}$$
$$T \rightarrow T \left[ 1 - \left( \frac{\omega}{\Omega_{\perp}} \right)^2 \right]^{1/3}$$

$$T_c = T_c^0 \left[ 1 - \left( \frac{\omega}{\Omega_{\perp}} \right)^2 \left\{ \frac{1}{3} + \frac{2\pi^2(\delta + 2\epsilon_F)}{3g^2\sqrt{\epsilon_F}} \left( \frac{\hbar^2}{2m} \right)^{3/2} + \frac{21}{2^{3/2}\pi^{7/2}} \frac{(\beta_c \epsilon_F)^2 \hbar \Omega_{\perp}}{\log[\beta_c \epsilon_F] \epsilon_F} \right\} \right]$$



# Non-interacting Phase Diagram

$$T_F(\omega) = T_F^0 \left(1 - \frac{\omega^2}{\Omega_{\perp}^2}\right)^{1/3}$$



# Symmetries of Hamiltonian

- Conservation of Particle Number

$$N = \int d^3\mathbf{r} \left( 2\langle \phi^\dagger(\mathbf{r})\phi(\mathbf{r}) \rangle + \sum_{\sigma} \langle \psi^\dagger_{\sigma}(\mathbf{r})\psi_{\sigma}(\mathbf{r}) \rangle \right)$$

- Angular Momentum  $[\mathcal{H}, L^z] = 0$
- Translation Symmetry in the x-y plane in the limit  $\omega = \Omega_{\perp}$

# Summary

- Investigate Critical Temperature of Fermion-Boson Model in a rotating trap
- Main effect of rotation
  - Lowering of  $T_c$  through a decrease of local density of atoms
- Future Work:
  - Derivation of G-L equation in rotating trap
  - Study  $\omega = \Omega_{\perp}$  regime